

§1 Introduction

Montag, 21. November 2022 13:50

§1.1 what are thin-films?

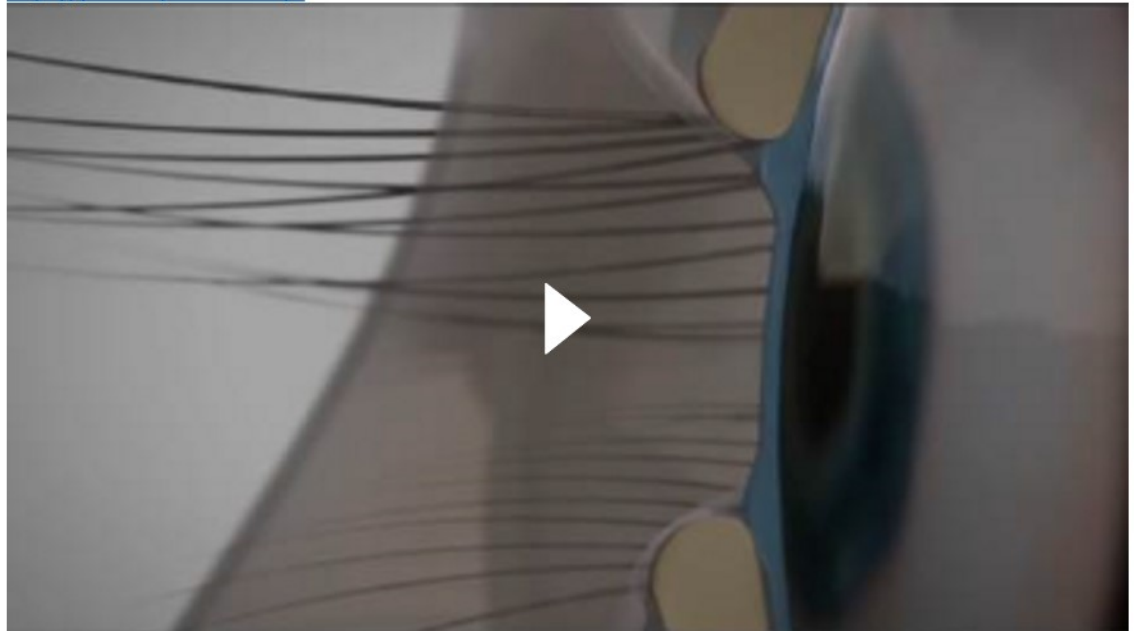
* a body of fluid, shallow in one dimension

Examples

- Tear film in the eye

* surface tension & gravity are at play

YT: AllThingsEyes
<https://youtu.be/WhxMfIG1Vpc>



- Lava flows:

- on large scales: a few meters are thin
- viscosity is key factor



YT: Olivier Grunewald [My drone above the incredible Icelandic geyser of lava](#)

- **Analogue Film production**
 - liquid-liquid interactions
 - don't want turbulence! want laminar flow.
 - liquid-solid interactions

YT: SmarterEveryDay
<https://youtu.be/cAAJUWh9F4>

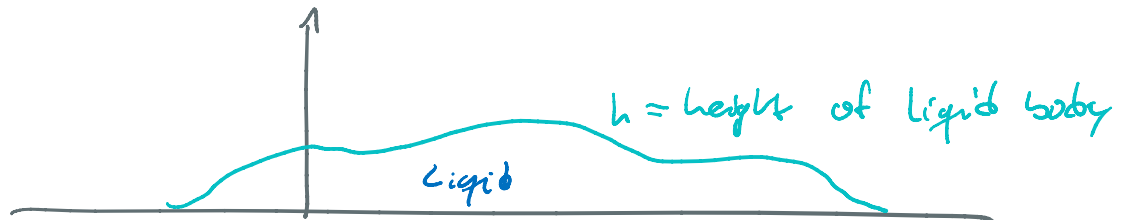


§1.2 what are thin-film equations?

• main idea: By assuming $\frac{\text{height}}{\lambda} \ll 1$

* Main idea: By assuming $\frac{\text{height}}{\text{length}} \ll 1$

we can reduce the question: "How does the liquid evolve?" to a single equation for the fluid's height.

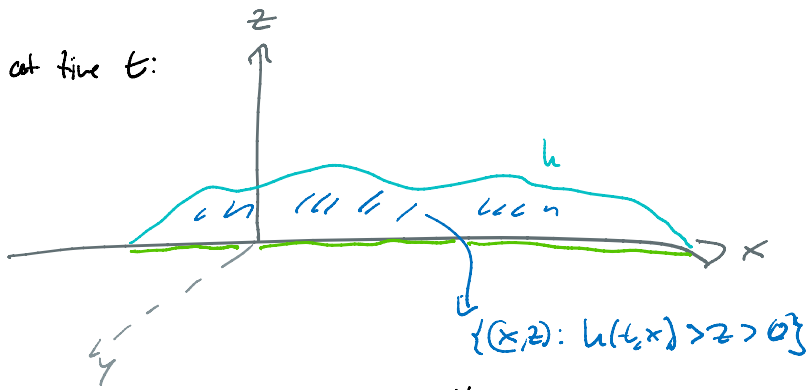


* Cannon in literature: only in one horizontal direction.

Goal derive this equation.

§2 Derivation of a thin-film equation

Mittwoch, 23. November 2022 10:01



need physics at three "zones"

- 1) Interior
- 2) Upper boundary
- 3) Lower boundary.

Zone 1 Interior

we use incompressible Navier-Stokes system $\mu a = F$

$u = u(t,x,z) \hat{=}$ horizontal velocity

$v = v(t,x,z) \hat{=}$ vertical velocity

$\pi = \pi(t,x,z) \hat{=}$ pressure

$\rho = \text{const} \hat{=}$ density

$\mu = \text{const} \hat{=}$ viscosity

$$\Rightarrow \begin{cases} \rho(\partial_t u + u \partial_x u + v \partial_z u) = -\partial_x \pi + \mu(\partial_x^2 + \partial_z^2)u \\ \rho(\partial_t v + u \partial_x v + v \partial_z v) = -\partial_z \pi + \mu(\partial_x^2 + \partial_z^2)v \\ \partial_x u + \partial_z v = 0 \end{cases}$$

inside $\{(t,x,z): h(t,x) > z > 0\}$

Zone 2 Upper boundary

• need to couple h to the system.

Kinematic equation: $\partial_t h + u \partial_x h = v$

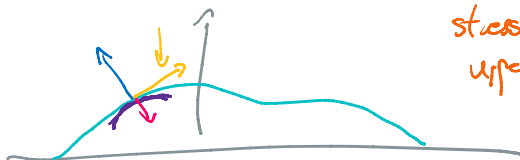
• Capillary equation:

$$\left(-\pi \mathbb{1} + \frac{\mu}{2} (D_x^2 u + D_x^2 v)^T \right) \begin{pmatrix} -\partial_x h \\ 1 \end{pmatrix} = \underbrace{\sigma}_{\text{surface tension}} \frac{\partial_x^2 h}{(1 + (\partial_x h)^2)^{3/2}} \begin{pmatrix} -\partial_x h \\ 1 \end{pmatrix}$$

stresses at the upper boundary

normal on the graph of h .

curvature of the graph of h



Zone 3 Lower boundary

Q: How do liquid particles interact with a solid?

• No-slip assumption: $u = 0$ at $z = 0$

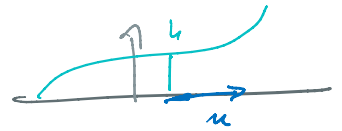
This has the no-slip paradox:

$\uparrow h$

• No-slip assumption: $u = v = 0$

This has the no-slip paradox!

• Navier-slip assumption: $u = 2 \partial_z u$ at $z = 0$



• More general assumption exist. $u = 2 \frac{3-n}{h^{n-2}} \partial_z u$ $n = 2$
 Also $v = 0$!
 $n \in (1, 3)$

§2.2 Lubrication approximation

- Navier-Stokes eqn
- Capillary equation
- Kinematic equation
- No-slip

look at this non-dimensionalised

• $x = L \bar{x}$, $z = H \bar{z}$, $\rho = \frac{\mu}{L}$, $t = \frac{\mu}{L^3} \bar{t}$

• $u = \frac{L}{\mu} \varepsilon^3 \bar{u}$, $v = \frac{H}{\mu} \varepsilon^3 \bar{v}$, $\pi = \frac{\mu}{L} \varepsilon \bar{\pi}$

\Rightarrow pressure \ll viscosity, \sim surface tension.

• $h = H \bar{h}$, $\sigma = \frac{\mu L}{\mu} \bar{\sigma}$, $Re = \frac{\rho \varepsilon^3 L^2}{\mu}$

$\Rightarrow \varepsilon^2 Re (\partial_x \bar{u} + \bar{u} \partial_x \bar{u} + \bar{v} \partial_z \bar{v}) = -\partial_x \bar{\pi} + (\varepsilon^2 \partial_x^2 \bar{u} + \partial_z^2 \bar{u})$

$\downarrow \varepsilon \rightarrow 0$

(1)
$$\begin{cases} \partial_x \bar{\pi} = \partial_z^2 \bar{u} \\ \partial_z \bar{\pi} = 0 \\ \partial_x \bar{u} + \partial_z \bar{v} = 0 \end{cases}$$

• Same goes for Kinematic equation $\rightarrow \partial_z \bar{h} + \bar{v} \partial_x \bar{h} = \bar{v}$

• Capillary eq: $\xrightarrow{\varepsilon \rightarrow 0}$ (2)
$$\begin{cases} \partial_z \bar{u} = 0 \\ \bar{\pi} = -\bar{v} \partial_x^2 \bar{h} \end{cases}$$
 at $\{\bar{h} = \bar{z}\}$

• $\bar{u} = 0, \bar{v} = 0$ (3)

Note: $Re = \frac{\rho \varepsilon^3 L^2}{\mu} \ll 1$. Small vertical velocities!
 \Rightarrow film need to stay small.

Goal compress this into a single equation for h .

Idea: express u, v, π in terms of $h \Rightarrow$ Kinematic equation.

usual compress ...

Idea express u, v, T in terms of $h \Rightarrow$ kinematic equation

At Π No more barred variables.

$$\Pi = -\sigma \partial_x^2 h \quad \text{at } z=h \quad \Rightarrow \quad \Pi = -\sigma \partial_x^2 h \quad \forall (x,z): h(t,x) > z > 0$$

$$\partial_t \Pi = 0$$

At u $\partial_z^2 u = \partial_x \Pi = -\sigma \partial_x^3 h$ in $h > z > 0$

$$\partial_z u = 0 \quad \text{at } z=h$$

$$u = 0 \quad \text{at } z=0$$

$$\partial_z u = \underbrace{\partial_z u|_{z=h}}_{=0} - \int_z^h \partial_z^2 u \, dz = \sigma \int_z^h \partial_x^3 h \, dz = \sigma (h-z) \partial_x^3 h$$

$$u = u|_{z=0} + \int_0^z \partial_z u \, dz = \sigma \int_0^z (h-\xi) \partial_x^3 h \, d\xi = \sigma \partial_x^3 h \left(z h - \frac{1}{2} z^2 \right)$$

At v similar calculation using $\partial_x u = -\partial_z v$ & $v|_{z=0} = 0$

At h insert u, v into $\partial_t h + u \partial_x h - v = 0$

$$0 = \partial_t h + u \partial_x h - v$$

$$= \partial_t h + \underbrace{\partial_x \left[\int_0^h u \, dz \right]}_{= \partial_x h \cdot u + \int_0^h \partial_x u \, dz} = \partial_x h \cdot u + \int_0^h \partial_x u \, dz$$
$$= \partial_x h \cdot u - \int_0^h \partial_z v \, dz$$

$$= \partial_t h + \partial_x \left[\int_0^h \sigma \partial_x^3 h \left(z h - \frac{1}{2} z^2 \right) dz \right] = \partial_x h \cdot u - v$$

$$= \partial_t h + \partial_x \left[\sigma \partial_x^3 h \left(\frac{1}{2} h^3 - \frac{1}{6} h^3 \right) \right]$$

$$= \partial_t h + \frac{\sigma}{3} \partial_x \left[h^3 \partial_x^3 h \right]$$

coeff of highest order derivative

• not 4-th order
• degenerate since the coeff. of the highest order can vanish.

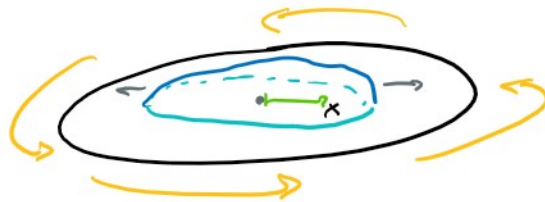
§3 Examples

Mittwoch, 23. November 2022 09:55

Examples

• Spin-coating (→ Guan et. al. 2017)

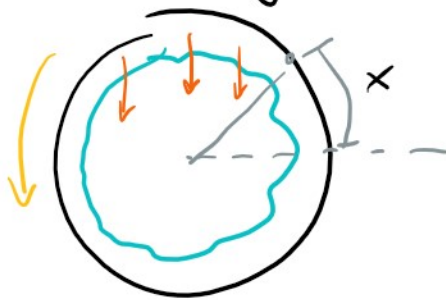
- * Fluid on a rotating plate
- * centrifugal forces act with surface tension



$$\Rightarrow \partial_t h + \underbrace{\frac{\partial_x [x^2 (h^2 + h^3)]}{x}}_{\text{centrifugal}} + \underbrace{\frac{\partial_x [x (h^2 + h^3) \partial_x \left[\frac{\partial_x [x \partial h]}{x} \right]]}{x}}_{\text{surface tension}} = 0$$

↳ 4th order equation.

• Rimming flows



→ Moffat 1977

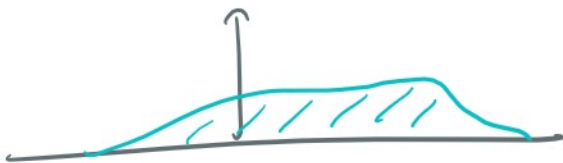
→ Rukhnachev 1977

• gravity acts

• outside rotates w/ constant speed.

$$\Rightarrow \partial_t h + \partial_x [h - b h^3 \cos(x) + a h^3 (\partial_x h + \partial_x^3 h)] = 0$$

• Horizontal plate w/o gravity



$$\Rightarrow \partial_t h + \partial_x [h^3 \partial_x^3 h] = 0.$$