

Guaranteed lift-off in non-Newtonian thin-film equations

March 7, 2023

Jonas Jansen

LTH, Lunds Universitet



LUNDS
UNIVERSITET

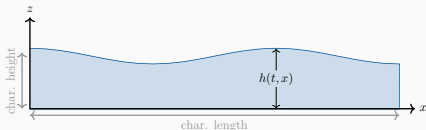
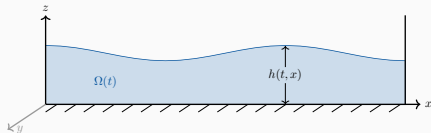
Non-Newtonian thin films

Thin fluid film

- incompressible
- viscous

$$\mu(\epsilon) = \mu_0 |\epsilon|^{\frac{1}{\alpha} - 1}$$

- capillary-driven



Navier-Stokes system: $\vec{u} = (u, v)$ velocity field in $\Omega(t)$

$$\left\{ \begin{array}{ll} \rho(\partial_t \vec{u} + (\vec{u} \cdot \nabla) \vec{u}), & = \operatorname{div} \cdot T(p, \vec{u}), \quad (t, x) \in \Omega(t) \\ \operatorname{div} \vec{u} & = 0, \quad (t, x) \in \Omega(t) \\ T(p, \vec{u}) \vec{n} & = \gamma \kappa \vec{n}, \quad z = h(t, x) \\ \partial_t h + u \partial_x h & = v, \quad z = h(t, x) \\ \vec{u} & = 0, \quad z = 0 \end{array} \right.$$

Properties

Thin-film equation

$$\begin{cases} \partial_t h + \partial_x (h^n |\partial_x^3 h|^{\alpha-1} \partial_x^3 h) = 0, & t > 0, x \in (0, 1) \\ \partial_x h = h^n |\partial_x^3 h|^{\alpha-1} \partial_x^3 h = 0, & t > 0, x \in \partial(0, 1) \\ h(0, x) = h_0(x), & x \in (0, 1) \end{cases}$$

- fourth-order doubly degenerate-parabolic equation
- mass conservation: $\int_{(0,1)} h(t, x) dx = \int_{(0,1)} h_0(x) dx$
- driven by surface tension: forces equilibrate surface area

$$E_s[h] = \int_{(0,1)} \sqrt{1 + |\partial_x h|^2} dx \cong 1 + \frac{1}{2} \int_{(0,1)} |\partial_x h|^2 dx = 1 + E[h]$$

- energy-dissipation mechanism

$$\frac{d}{dt} E[h](t) = - \int_{(0,1)} h^n |\partial_x^3 h|^{\alpha+1} dx \leq 0$$

Weak solutions

Thin-film equation

$$\begin{cases} \partial_t h + \partial_x (h^n |\partial_x^3 h|^{\alpha-1} \partial_x^3 h) = 0, & t > 0, x \in (0, 1) \\ \partial_x h = h^n |\partial_x^3 h|^{\alpha-1} \partial_x^3 h = 0, & t > 0, x \in \partial(0, 1) \\ h(0, x) = h_0(x), & x \in (0, 1) \end{cases}$$

Definition

$h \in L_\infty([0, \infty); H^1(\Omega)) \cap C^{\frac{1}{5\alpha+3}, \frac{1}{2}}([0, \infty] \times \bar{\Omega})$ with $\partial_x^3 h \in L_{\alpha+1, \text{loc}}(\{u > 0\})$ and $\partial_t h \in L_{\alpha+1}([0, \infty); (W_{\alpha+1}^1(\Omega))')$ is a weak solution to the thin-film equation if

$$\int_0^\infty \langle \partial_t h, \phi \rangle_{W_{\alpha+1}^1} dt - \iint_{\{u>0\}} h^n |\partial_x^3 h|^{\alpha-1} \partial_x^3 h \partial_x \phi dx dt = 0$$

for all $\phi \in L_{\alpha+1}((0, \infty); W_{\alpha+1, B}^1(\Omega))$.

Non-negativity of solutions

Question

Do solutions remain non-negative?

- For $\partial_t u + \Delta^2 u = 0$, the answer is no!
- If $\alpha = 1$ and $n > 1$, use entropy method:

$$\int_{(0,1)} g(h(t, x)) dx + \int_0^t \int_{(0,1)} |\partial_x^2 h| dx dt \leq \int_{(0,1)} g(h(0, x)) dx,$$

where $g''(h) = \frac{1}{h^n}$, hence $g(h) \sim h^{2-n}$ and $g(h) = +\infty$, $h < 0$

- For non-Newtonian rheologies: regularisation techniques

Existence of weak solutions

Theorem

Let $h_0 \in H^1((0,1))$, $h_0 \geq 0$. Then there exists a weak solution h to the thin-film equation that satisfies

$$E[h](t) + \int_0^t \int_{(0,1)} h^n |\partial_x^3 h|^{\alpha+1} dx ds \leq E[h_0] \quad (\text{EDI})$$

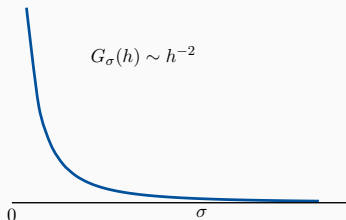
If $h_0 > 0$, then (EDI) is an equality as long as $h(t, \cdot) > 0$.

- Positivity by singular potential:

$$\partial_x^3 h \rightarrow \partial_x^3 h - G'_\sigma(h) \partial_x h$$

$$E[h] \rightarrow E[h] + \int G_\sigma(h) dx$$

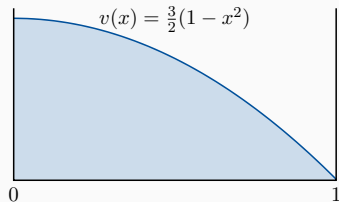
- Gradient-flow approach in $H^1(\Omega)$
- energy methods to identify limit as $\sigma \rightarrow 0$



The lift-off property

- $v(x) = \frac{3}{2}(1 - x^2)$ is steady state

$$\int_{(0,1)} v(x) dx = 1$$



- Observation

$$v = \arg \min E[h]$$

among all $h \geq 0$ with $\int h dx = \int v dx = 1$ and at least one root

- let $E[h_0] < E[v]$ and $\int h_0 dx = 1 \Rightarrow E[h](t) < E[v]$ for all t
- Conclusion: h solution with initial value h_0 , then $h(t) > 0$ and

$$E[h](t) + \int_s^t \int_{(0,1)} h^n |\partial_x^3 h|^{\alpha+1} dx d\tau = E[h](s)$$

The lift-off property

Question

Is the dissipation $\int_{(0,1)} h^n |\partial_x^3 h|^{\alpha+1} dx$ of order one?

Theorem

Let $n, \alpha > 0$ with $2(\alpha + 1) > n$. Then there exists $t_0 = t_0(n, \alpha) > 0$ such that whenever $h_0 \in H^1(\Omega)$ with $\int_{(0,1)} h_0 dx = 1$ and $E[h_0] < E[v]$, then any weak solution h to the thin-film equation with initial value h_0 satisfies

$$\min_{x \in \bar{\Omega}} h(t, x) \geq \frac{1}{2}, \quad t \geq t_0.$$

Long-time behaviour: convergence to mean with explicit convergence rates

Idea of proof

Lemma

Let $n, \alpha > 0$. There is a constant $C = C(n, \alpha) > 0$ such that for any $h \in W_{\alpha+1}^3(\Omega)$ with $h'(0) = h'(1) = 0$, $\int h \, dx = 1$ and $\min_{x \in \bar{\Omega}} h = \delta \in (0, 1/2)$ we have

$$\int_{\Omega} h^n |\partial_x^3 h|^{\alpha+1} \, dx \geq \frac{C}{\log^{\alpha+1}(M/\delta)} \min\{\delta^{n-1-2\alpha}, 1\}.$$

- $\min h = \delta < \frac{1}{2} \Rightarrow \max h > 1 \Rightarrow \exists x_0 : |h'(x_0)| > \frac{1}{4}$ and $h''(x_0) = 0$
- from $h'(1) = 0$, we conclude there is x_1 with $|h''(x_1)| = 1/2$
- By Hölder's inequality

$$\begin{aligned} \left(h''(x_0) + \frac{1}{2}\right)^{\alpha+1} &= \left(\int_{x_0}^{x_1} -\partial_x^3 h \, dx\right)^{\alpha+1} \\ &\leq \left(\int_{x_0}^{x_1} h^n |\partial_x^3 h|^{\alpha+1} \, dx\right) \left(\int_{x_0}^{x_1} h^{-\frac{n}{\alpha}} \, dx\right)^{\alpha} \end{aligned}$$

Idea of proof

Statement of Theorem

Want: $\min_{x \in \bar{\Omega}} h(t, x) \geq \frac{1}{2}, \quad t \geq t_0$

- As long as $\min_{x \in \bar{\Omega}} h(t, x) \leq \frac{1}{2}$:

$$\min_{x \in \bar{\Omega}} h(t, x) \geq C(E[v] - E[h](t))$$

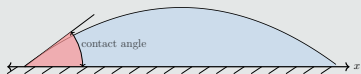
$$\begin{aligned} E[v] - E[h](t) &> C(E[h_0] - E[h](t)) \\ &= C \int_0^t \int_{(0,1)} h^n |\partial_x^3 h|^{\alpha+1} dx ds \\ &\geq C(\min_{x \in \bar{\Omega}} h(t, x))^{1-\varepsilon} \end{aligned}$$

- Conclusion: $\frac{d}{dt}(E[v] - E[h](t)) \geq C(E[v] - E[h](t))^{1-\varepsilon}$
- Gronwall gives: $E[h](t) \rightarrow 0$ if $\min h(t, x) < \frac{1}{2}$ for all times

Further projects and questions

Droplets

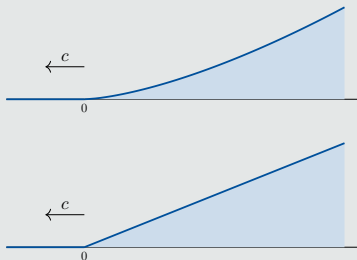
Study of droplet case



Question: Find solutions to full free-boundary problem

Travelling waves

Behaviour near contact points:



$h(t, x) = H(x - ct)$ and obtain ODE

$$cH = H^n |H'''|^{\alpha-1} H'''$$

Thank you for your attention!

Questions?